

Thin film flow of non-Newtonian fluids on a vertical moving belt using Homotopy Analysis Method

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Abstract

In this study, the problem of thin film flow of two non-Newtonian fluids namely, a Sisko fluid and an Oldroyd 6-constant fluid on a vertical moving belt that modeled by a system of nonlinear differential equations is studied. The system is solved using the Homotopy Analysis Method (HAM), which yields an analytic solution in the form of a rapidly convergent infinite series with easily computable terms. Homotopy analysis method contains the auxiliary parameter \hbar , which provides us with a simple way to adjust and control the convergence region of solution series. By suitable choice of the auxiliary parameter \hbar , reasonable solutions for large modulus can be obtained.

Keywords: thin film, non-Newtonian fluids, HAM.

Introduction

Recently, finding analytical approximating solutions of nonlinear equations has widespread applications in numerical mathematics and applied mathematics and there has appeared an ever increasing interest of scientists and engineers in analytical techniques for studying nonlinear problems. Homotopy Analysis Method has been proposed by Liao and a systematic and clear exposition on this method is given in [1-2]. HAM is a powerful mathematical technique and has already been applied to several nonlinear problems [3-21]. This method contains an auxiliary parameter h which provides us with a simple way to adjust and control the convergence region of the solutions. The equations modeling non-Newtonian incompressible fluid flow give rise to highly nonlinear differential equations. Such non-Newtonian fluids find wide applications in commerce, industry and have now become the focus of extensive study.

In this study, HAM is applied to problem of thin film flow of two non-Newtonian fluids namely, a Sisko fluid [22] and an Oldroyd 6-constant fluid on a vertical moving belt. Using HAM, we have obtained meaningful solutions for ψ which clearly shows the capability of this method for solving equations with high nonlinearities and controlling the convergence region of solutions. The HAM solutions for two considered examples in comparison with HPM solutions of [23] admit a remarkable accuracy.

1. Basic ideas of homotopy analysis method

To describe the basic ideas of the HAM, consider the following differential equation,

$$\mathcal{N}[y(t)] = 0, \quad (1)$$

where \mathcal{N} is a nonlinear operator, t denotes the independent variable, and $y(t)$ is an unknown function respectively. By means of generalizing the traditional homotopy method, we have the so-called zeroth-order deformation equation as

$$(1-q)\mathcal{L}[\psi(t,q) - y_0(t)] = q\hbar\mathcal{N}[\psi(t,q),q], \quad (2)$$

where $\psi(t,q)$ is an unknown function, \mathcal{L} is an auxiliary linear operator, $q \in [0,1]$ an embedding parameter, \hbar a non-zero auxiliary parameter, and $y_0(t)$ is an initial guess of $y(t)$. It is important to note that, one has freedom to choose auxiliary parameters such as \hbar and \mathcal{L} in HAM. Obviously, when $q=0$ and $q=1$, both

$$\psi(t;0) = y_0(t) \quad \text{and} \quad \psi(t;1) = y(t). \quad (3)$$

Thus as q increases from 0 to 1, the solution $\psi(t,q)$ varies smoothly from the known initial guess $y_0(t)$ to the solution $y(t)$. By expanding $\psi(t,q)$ into Taylor series of the embedding parameter q and using (3), we have

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$$\psi(t; q) = y_0(t) + \sum_{n=1}^{+\infty} y_n(t) q^n, \quad (4)$$

where

$$y_n(t) = \frac{1}{n!} \left. \frac{\partial^n \psi(t; q)}{\partial q^n} \right|_{q=0}, \quad (5)$$

If the auxiliary linear operator, the initial guess, the auxiliary parameter \hbar , are so properly chosen, then the series (4) converges at $q=1$ and

$$y(t) = y_0(t) + \sum_{n=1}^{+\infty} y_n(t) \quad (6)$$

The governing equations of y_n can be deduced from the zeroth-order deformation equations (2) and (3). Define the vector

$$\vec{y}_n = \{y_0(t), y_1(t), \dots, y_n(t)\}, \quad (7)$$

Now substituting Eq. (5) into Eq. (2) and differentiating Eq. (2) n times with respect to the embedding parameter q , then dividing by $n!$, and finally setting $q=0$, we have the n^{th} -order deformation equation

$$\mathcal{L}[y_n(t) - \chi_n y_{n-1}(t)] = \hbar R_n(y_{n-1}), \quad (8)$$

subjected to boundary conditions

$$y_n(0) = 0, \quad y'_n(0) = 0, \quad (9)$$

where prime denotes the derivative with respect to t and

$$R_n(y_{n-1}) = \frac{1}{(n-1)!} \left. \frac{\partial^{n-1} \mathcal{N}[\psi(t; q), q]}{\partial q^{n-1}} \right|_{q=0}, \quad (10)$$

and

$$\chi_k = \begin{cases} 0 & k \leq 1 \\ 1 & k > 1 \end{cases} \quad (11)$$

It should be noted that $y_n(t)$ ($n \geq 1$) is governed by the linear equation (8) together with the linear boundary conditions that come from the original problem.

Problem formulation

We consider a container having a non-Newtonian fluid in it [24-29]. A wide moving belt passes through this container, which moves vertically upward with constant velocity U_0 as shown in Fig.1 Since the belt moves upward and passes through the fluid, it picks up a film fluid of thickness d . Due to gravity, the fluid film

tends to drain down the belt. For simplicity, the following assumptions are made:

- (i) The flow is in steady state.
- (ii) The flow is laminar and uniform.
- (iii) The film fluid thickness d is uniform.

We choose an xy -coordinate system and position x -axis parallel to the fluid and normal to the belt, y -axis upward along the belt and z -axis normal to the xy -plane. The only velocity component is in the y -direction, therefore,

$$V = (0, v(x), 0) \quad (12)$$

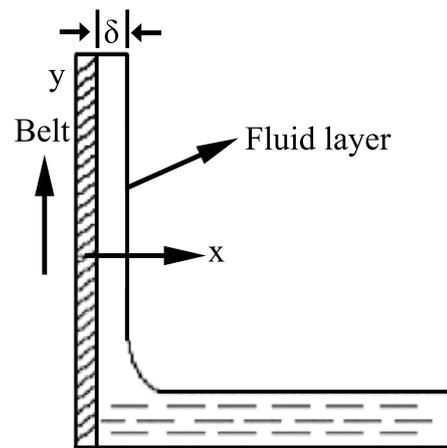


Figure 1. Geometry of the flow of moving belt through a non-Newtonian fluid

and also the extra stress tensor is function of x only, that is

$$S = S(x), \quad (13)$$

Eq. (12) satisfies the continuity Equation identically.

In the sequel, we first derive the flow equations for a Sisko fluid and an Oldroyd 6-constant fluid separately and then solve the resulting equations by using homotopy perturbation method.

Governing equation for a Sisko fluid

Consider the following differential equation of a Sisko fluid over a moving belt [30]

$$\frac{d^2 v}{dx^2} + nb \left[\frac{dv}{dx} \right]^{n-1} \frac{d^2 v}{dx^2} - k = 0 \quad (14)$$

subjected to the boundary conditions

$$\begin{aligned} v &= 1 & x &= 0, \\ \frac{dv}{dx} &= 0 & x &= 1, \end{aligned} \quad (15)$$

For HAM solutions, we choose the initial guesses and auxiliary linear operator in the following forms:

$$v_0'' = 0, \rightarrow v_0 = \frac{k}{2}(x^2 - 2x) - k, \quad (16)$$

Zeroth –order deformation problems

$$(1-p)L[v(x, p) - v_0(x)] = p\hbar N[v(x, p)], \quad (17)$$

$$v(0, p) = 0 \quad v'(0, p) = 1 \quad (18)$$

$$N[v(x, p)] = \frac{d^2v(x, p)}{dx^2} + nb \left[\frac{dv(x, p)}{dx} \right]^{n-1} \frac{d^2v(x, p)}{dx^2} - k, b = const. \quad (19)$$

For $p=0$ and $p=1$ we have:

$$v(x, 0) = v_0(x), \quad v(x, 1) = v(x),$$

When p increases from 0 to 1 and $v(x, p)$ vary from $v_0(x)$ to $v(x)$. Due to Taylor’s series with respect to p we have:

$$v(x, p) = v_0(x) + \sum_{m=1}^{\infty} v_m(x) p^m, \quad (20)$$

$$v_m(x) = \frac{1}{m!} \frac{\partial^m (v(x, p))}{\partial p^m}$$

In which \hbar is chosen in such a way that this series is convergent at ($p=1$).

Therefore we have through Eq. (20) that:

$$v(x) = v_0(x) + \sum_{m=1}^{\infty} v_m(x), \quad (21)$$

Convergence of the HAM solution

As pointed out by Liao, the convergence and rate of approximation for the HAM solution strongly depends on the values of auxiliary parameter \hbar . Fig. (2) clearly depict that the ranges, for admissible values of \hbar is $-1.4 < \hbar < -0.6$. Our calculations clearly indicate that series (21) converge for whole region of x when $\hbar = -1$.

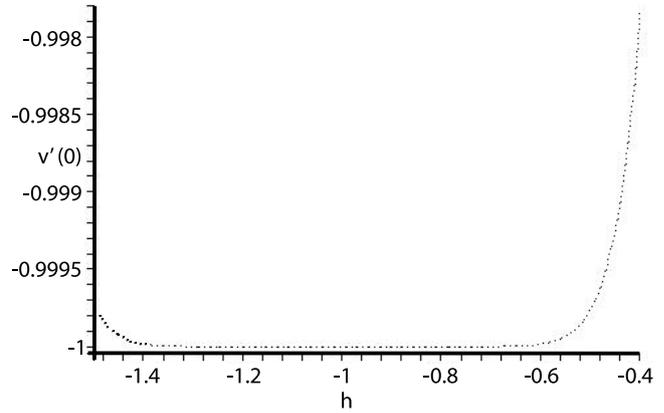


Figure 2. Variation of \hbar for 12th-order of approximation

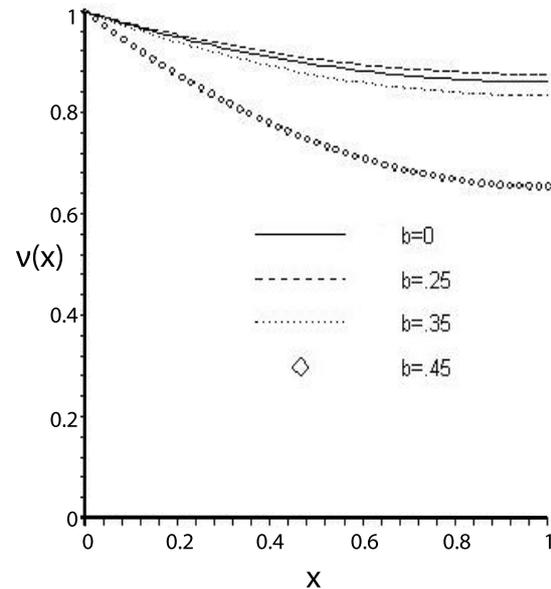


Figure 3. Variation of $v(x)$ for $n=1$ and $k=1$

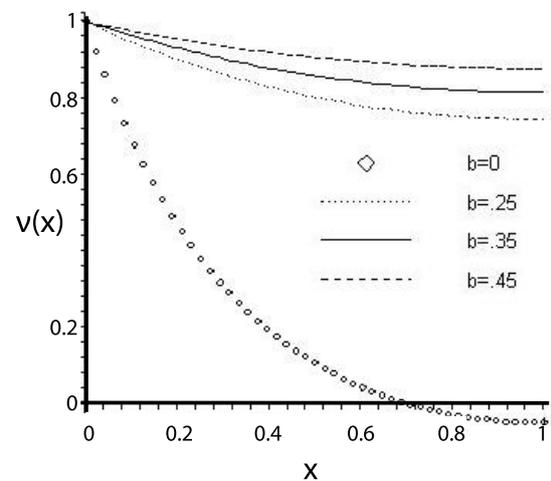


Figure 4. Variation of $v(x)$ for $n=2$ and $k=1$

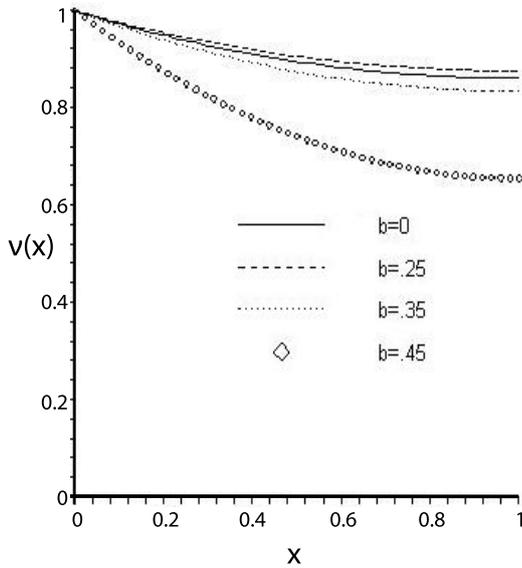


Figure 5. Variation of $v(x)$ for $n=3$ and $k=1$

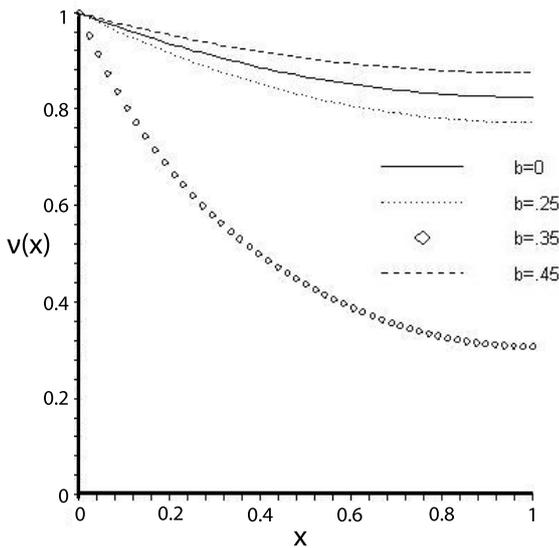


Figure 6. Variation of $v(x)$ for $n=4$ and $k=1$

Application of HAM to the problem of Oldroyd 6-constant fluid

Our differential equation with the boundary conditions is [30]

$$\frac{d^2v}{dx^2} + (3a_1 - a_2)\left(\frac{dv}{dx}\right)^2 \frac{d^2v}{dx^2} + a_1a_2\left(\frac{dv}{dx}\right)^4 \frac{d^2v}{dx^2} - m\left(1 + a_2\left(\frac{dv}{dx}\right)^2\right)^2 = 0, a_1, a_2 = const. \tag{22}$$

$$v = 1 \quad x = 0, \\ \frac{dv}{dx} = 0 \quad x = 1, \tag{23}$$

For HAM solutions, we choose the initial guesses and auxiliary linear operator in the following forms:

$$v_0'' = 0, \rightarrow v_0 = \frac{m}{2}(x^2 - 2x) - m, \tag{24}$$

$$(1-p)L[v(x,p) - v_0(x)] = p\hbar N[v(x,p)], \tag{25}$$

$$v(0,p) = 0 \quad v'(0,p) = 1$$

$$N[v(x,p)] = \frac{d^2v}{dx^2} + (3a_1 - a_2)\left(\frac{dv}{dx}\right)^2 \frac{d^2v}{dx^2} + a_1a_2\left(\frac{dv}{dx}\right)^4 \frac{d^2v}{dx^2} - m\left(1 + a_2\left(\frac{dv}{dx}\right)^2\right)^2, \tag{26}$$

For $p=0$ and $p=1$ we have:

When p increases from 0 to 1 $v(x,p)$ vary from $v_0(x)$ to $v(x)$. Due to Taylor's series with respect to p we have:

$$v(x,p) = v_0(x) + \sum_{m=1}^{\infty} v_m(x)p^m, \tag{27}$$

$$v_m(x) = \frac{1}{m!} \frac{\partial^m(v(x,p))}{\partial p^m}$$

In which \hbar is chosen in such a way that this series is convergent at ($p=1$).

Therefore we have through Eq. (27) that:

$$v(x) = v_0(x) + \sum_{m=1}^{\infty} v_m(x), \tag{28}$$

Convergence of the HAM solution

As pointed out by Liao, the convergence and rate of approximation for the HAM solution strongly depends on the values of auxiliary parameter \hbar . Fig.(8) clearly depict that the ranges, for ad-

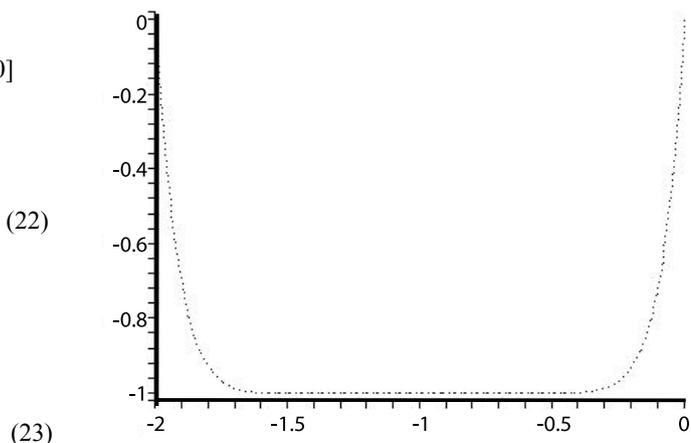


Figure 7. Variation of \hbar for 8th-order of approximation

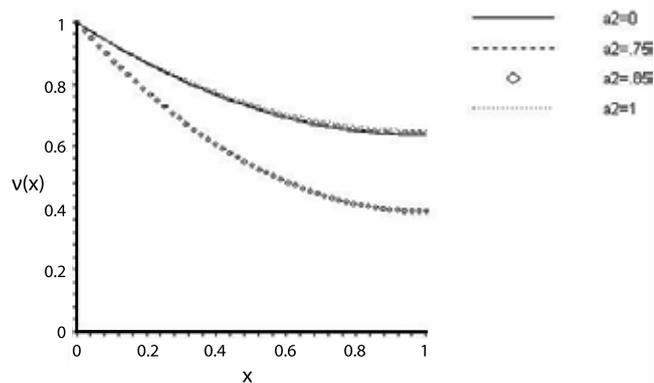


Figure 8. Variation of $v(x)$, for $m=1$ and $\alpha_1=0.1$

missible values of \tilde{h} is $-1.7 < \tilde{h} < -0.3$. Our calculations clearly indicate that series (28) converge for whole region of x when $\tilde{h} = -1$.

Conclusions

In this paper, the capability of HAM for obtaining approximate solutions of the velocity profile of thin film flow of non-Newtonian fluids over a moving belt is shown. First HAM is applied to the problem of a Sisko fluid and the problem of an Oldroyd 6-constant fluid is considered for further studies. Using HAM, we have achieved an analytical solution for the velocity of the belt. The results are compared with those of obtained via HPM which clarifies the effectiveness of this method. It is observed from the figures that the speed of the belt decreases as the non-Newtonian effects increase.

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