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The Creep Properties of Fine Sandstone under Uniaxial Tensile Stress

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Abstract

A graduated uniaxial direct tensile creep test for fine sandstone is conducted by adopting a custom-designed direct tensile test device for rock. The experiment shows that the tensile creep of fine sandstone has similar creep curve patterns to those of compression creep, while the ratios of the creep strain to the total strain obtained in the tensile tests are substantially higher than those obtained for similar compression tests, which indicates that the creep ability of rock in the tensile process is higher than that in the uniaxial compression process. Based on the elastic modulus in the approximately linear portion of the obtained isochronous stress-strain curves of the tensile creep, the time dependence of the elasticity modulus for the Kelvin model is evaluated, and a revised generalized Kelvin model is obtained by substitution into the generalized Kelvin model. A new viscousplastic model is proposed to describe the accelerated creep properties, and this model is combined in series with the revised generalized Kelvin model to form a new nonlinear viscoelastic-plastic creep model that can describe the properties of attenuation creep, steady creep, and accelerated creep. Comparison of the test and theoretical curves demonstrates that they are nearly identical, which verifies the performance of the model.

Keywords: Fine sandstone, Direct tensile test, Creep, Nonlinear creep model, Model identification

1. Introduction

In the traditional view, if it is assumed that rock can sustain tensile stress at all, the deformations of rock under both compression and tension loads are generally identified as equivalent. However, in actuality, the uniaxial compressive strength of rock is much higher than its tensile strength, with the ratio of the two strengths ranging from 6 to 80 [1]. The difference between the compressive and tensile strengths of rock is so great, in fact, that actual rock engineering requirements obviously cannot be met under the assumption that deformations owing to the two stresses are identical. In the construction excavation process of slope engineering, mining engineering, and underground engineering, tensile stress will always be produced [2]. Moreover, tensile stresses may by no means lead to instability immediately after rock excavation and exploitation, but, rather, may require a considerable period of time. Because rock has little tensile strength, the presence of tensile stresses in rock engineering projects can be a decisive factor for the engineering stability. Therefore, the study of the creep properties of rock under conditions of tensile stress is of theoretical and practical significance.

Currently, thorough studies of the creep properties of rock under conditions of compressive stress are abundant and have achieved a great deal. Mishra et al. investigated the time-dependent behavior of laminated shale based on the triaxial creep tests [3]. Gabriel et al. proposed a new empirical model that comprehends the three creep stages [4]. Wassmann et al. investigated the rheology of the plate interface dissolution precipitation creep in high pressure metamorphic rocks [5]. Julie et al. studied how interactions between brittle and viscous mechanisms lead to widespread transformation of the rocks and how a shear zone may evolve from a zone prone to earthquakes and postseismic creep to a zone of steady state creep [6]. Wei et al. developed a grain-based heterogeneous numerical model originated from the discrete element method, and the procedure is successfully applied to simulate uniaxial compression tests, Brazilian tests and uniaxial creep tests [7]. Gu et al. investigated the creep properties of sandstone under uniaxial compression [8]. M Schoenball et al. proposed to explain observations of delayed wellbore failure by timedependent brittle creep, which has been observed for many types of rocks [9]. Siavash et al. investigate the timeindependent behavior of grouting material and the creep tests of grouting material was conducted in triaxial compression apparatus at room temperature [10]. Hakan et al. proposed a mathematical modeling of the creep behavior of Tuzkoy rock salt based on the creep tests [11]. Karimpour et al. conducted triaxial compression tests on dense specimens of Virginia Beach sand at low and high confining pressures to study time effects that relate to grain crushing due to static fatigue or delayed fracture [12]. However, studies of the creep properties of rock under conditions of direct tensile stress are rare.

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The main reason for the rareness in direct tensile creep properties are two-fold. Firstly, the restrictions of test conditions inhibit study in a laboratory setting, and, secondly, actual rock engineering projects predominantly consider conditions of compressive stress. However, with increasing investment in infrastructure, numerous large-scale and structurally complex rock engineering projects are presently under construction, which involve an increasing number of issues concerning rock deformation and instability under conditions of tensile stress. Therefore, the author has designed a direct tensile device for rock that can better avoid eccentric tension. Fine sandstone from a foundation pit in Chongqing was selected as the study object, and graduated uniaxial direct tensile creep testing was conducted. By considering the time dependence of the elastic modulus in the approximately linear portion of the obtained isochronous stress-strain curves, the equation for the viscoelastic modulus of a Kelvin model with respect to time is formulated, and a revised generalized Kelvin model is obtained by substitution into the generalized Kelvin model. The model obtained is combined with a newly proposed viscoplastic model in series to form a new viscoelasticplastic creep model. The model parameters are identified utilizing the 1stOpt mathematics optimization software. The results presented in the present study may be of significant value to the stability issues of rock engineering.

2. Creep test

2.1 Test equipment

The test adopts a custom-designed direct tensile device for rock, which is mainly composed of a longitudinally-symmetric force support element (namely, a pull rod) and a pull head as well as high-strength nut, as shown in Figure 1. The test incorporates the simple loading device shown in Figure 2. The working principle is described in detail elsewhere [2].



Fig. 1. Rock uniaxial direct tensile device. (a) schematic diagram. (b) image of actual device



Fig. 2. Loading device for uniaxial direct tensile creep test. (a) schematic diagram . (b) image of actual device

2.2 Sample preparation and testing program

The rock sample used for testing is a fine sandstone obtained from a foundation pit in Chongqing, which is mainly comprised of quartz, feldspar, flint, and muscovite. The dry density is 2.33–2.74 g/cm3, porosity 0.63–0.71%, grain diameter 0.01–0.5 mm, uniaxial compressive strength 63.40 MPa, elasticity modulus 25.03 GPa, and Poisson ratio 0.13. Cylindrical samples of 50 mm diameter and 100 mm height were processed in accordance with ISRM testing regulations.

As shown in Figure 1, an axial and a lateral strain gages were pasted symmetrically in the middle region of each sample, which are used to measure the axial and radial strain of the sample. A DH3818 static strain test system comprised of a data collection box, microcomputer, and supporting software was adopted to collect the strain and stress of the specimen during the loading process, and is able to quickly measure the static strain and stress values of several points in large structures, models, and materials stress tests automatically, accurately, and reliably. Prior to conducting the tensile creep tests, tensile strength tests were first conducted on three fine sandstone samples, and the average tensile strength and tensile yield stress values were determined to be 1.78 MPa and 1.25 MPa, respectively. The reason for that is to determine the stress level in tensile creep test.

The creep tests were divided into two groups, and Table 1 lists the graduated test loading programs employed for each.

	Table 1. Creep	test programs for	graduated loading	of uniaxial tension.
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Specimen	Stress	Stress level (MPa) /stress-strength ratios				
No.	Level I	Level II	Level III	Level IV		
Ι	0.44/0.25	0.74 /0.40	1.04 /0.58	1.34 /0.75		
II	0.53 /0.30	0.84 /0.47	1.14 /0.64	1.45 /0.81		

2.3 Creep deformation analysis

Figure 3 presents the tensile creep curves of the two specimens under the loading programs listed in Table 1, where the various levels of stress loading are identified above the curve.



Fig. 3. Uniaxial tensile creep curves for fine sandstone. (a) specimen I. (b) specimen II

Based on the test results, the axial creep strain and total strain of the rock under the various levels of stress are listed in Table 2.

 Table 2. Axial tensile creep strain and total strain of fine sandstone under the graduated loading listed in table 1.

Specimen No.	Stress (MPa)/	Axial strain			
	stress-strength ratios	Creep strain	Total strain	Ratio of creep strain to total strain (%)	
	0.44 /0.25	0.3442	1.5327	22.45	
Ι	0.74 /0.40	0.6475	1.3115	49.37	
	1.04 /0.58	0.9385	1.7008	55.18	
	1.34 /0.75	1.0451	1.8648	56.04	
Π	0.53 /0.30	0.7861	2.2042	35.66	
	0.84 /0.47	1.2909	2.0655	62.49	
	1.14 /0.64	1.3294	1.8689	71.13	
	1.45 /0.81	0.6898	1.2794	53.91	

It can be seen from Figure 3 that, for stress levels less than the yield stress, the creep curve can be divided into two stages: an attenuation creep stage and a steady creep stage. Under the highest stress level, a comparatively complete creep curve with three stages appear, that is, an attenuation creep stage, steady creep stage, and accelerated creep stage. The creep strain value increases with increasing stress, and the ratio of the creep value to the total strain also increases, which indicates that the creep strain of the rock accounts for an increasing percentage of the total deformation with increasing stress.

The processes of uniaxial compression creep shown in Figure 4, present a similar behavior to those processes shown in Figure 3. However, by comparing the test results of uniaxial tension and compression, we can find that the creep processes of the two differ greatly. In the tensile creep test, the stress-strength ratios of specimens I and II are 0.25 and 0.30, respectively, and the corresponding ratios of the creep strain to the total strain are respectively 22.45% and 35.66%, while those reach 55.18% and 62.49%, respectively, when the stress-strength ratios are respectively 0.58 and 0.47.

However, in the literature [8], for the uniaxial compressive creep test of No. 4 fine sandstone, when the stress-strength ratio is 0.25, the ratio of the creep strain to the total strain is only 7.21%. In addition, when the stress-strength ratio of NO.3 fine sandstone is 0.30, the corresponding ratio of the creep strain to the total strain is 4.35% [8]. It can be observed that, under the same or similar stress-strength ratios, the ratio of the creep strain to total strain in the uniaxial tensile creep process is considerably higher than the corresponding value in the uniaxial compressive creep process.



Fig. 4. Uniaxial compression creep curve for fine sandstone [8] (which belong to the same grade of sandtone with the specimen in main text). (a)uniaxial compression creep curve for a No. 3 specimen. (b) uniaxial compression creep curve for a No. 4 Specimen

2.4 Isochronous stress-strain curve properties

Presently, studies of the isochronous stress-strain curves for rock under compressive stress are abundant, whereas those studies for rock under tensile stress are rare. Here, further studies are conducted of the isochronous stress-strain curves for rock under tensile stress. To this end, we select the strain values of the creep curves of specimens I and II at regular points of time under the various stress levels and obtain the stress-strain isochronous curves shown in Figure 5 for the fine sandstone samples during the tensile creep process.





Fig. 5. Isochronous stress-strain curves. (a) specimen I. (b) specimen II

As observed in Figure 5, at lower stress levels, the creep deformation of the rock and the stress are essentially in a linear relationship. However, with increasing stress level, the isochronous stress-strain curve appears to be of an obvious nonlinear relationship. In addition, with increasing time, the isochronous stress-strain curves gradually deviate toward the strain axis, which indicates that the viscous deformation properties of the rock sample tend to develop with increasing strain.

3. Methodology

3.1 Revised generalized Kelvin model and its creep equation

From the results of uniaxial tensile test studies in recent years [13], [14], [15], [16], we find that the general shape of the stress-strain curves for both uniaxial tension and uniaxial compression are similar. Therefore, in reference to a creep model under compressive stress, a constitutive model to describe rock creep under tensile stress can be developed.

Figure 6 illustrates the Kelvin model in which the stress and strain of the spring body have a linear relationship, which means that the elasticity modulus is a constant. The stress-strain curves for most rock materials demonstrate such a relationship, while, in such a model, because the spring body and the viscous device are in parallel, the stress and strain of the spring body are no longer in a simple linear relationship. In the creep process, spring body deformation develops gradually with the passage of time, and ends instantly. Therefore, in the creep process, the elasticity modulus E of the spring body is no longer a constant, and changes continuously with the passage of time, such that E is given as E(t). The nature of the time dependence of E(t) can be obtained from the isochronous stress-strain curves of the creep process.



In Figure 6, when the stress level is less than the yield stress, the rock is mainly a viscoelastic material, corresponding to the approximately linear portions of the isochronous stress-strain curves. Therefore, a curve of the elasticity modulus with respect to time can be obtained through determining the elasticity modulus of the approximately linear portion at different points on the isochronous stress-strain curve, as shown in Figure 7.



Fig. 7. Elasticity modulus as a function of time.(a)specimen I. (b) specimen II

The test data of the elasticity moduli with respect to time shown in Figure 7 were fit using the form $E(t) = \alpha \exp(-\beta t)$ where α is a fitting parameter and β is a positive constant representing the material nonlinearity, and the results listed in Table 3. From the R² values given in Table 3, we can observe that the description of the creep viscoelastic modulus with respect to time using an exponential function is suitable.

Table 3. Creep test programs for graduated loading of uniaxial tension fitting results of $E(t) = \alpha \exp(-\beta t)$.

	Fitting parameters				
Specimen No.	α	β	R^2		
Ι	2.5015	0.0052	0.9891		
II	1.5018	0.0010	0.9377		

If it is assumed that the elasticity modulus is E_0 at t = 0, then the boundary condition $E(t) = E_0$ is established, and the following form of E(t) can be obtained.

$$E(t) = E_0 \exp(-\beta t) \tag{1}$$

A revised generalized Kelvin model is thus obtained by assuming that the elasticity modulus E_2 of the spring body of the Kelvin model fit the function given by Eq.1, as shown in Figure 8, and assuming the elasticity modulus of the other spring is E_1 .



Fig. 8. Revised generalized Kelvin model.

The corresponding state equation of the revised generalized Kelvin model is as follows:

$$\sigma = E_1 \varepsilon_1 = E_2 \varepsilon_2 + \eta \varepsilon_2$$

$$E_2 = E_0 e^{-\beta t}$$
(2)

Here, σ is the stress, ε is the strain, and η the viscosity associated with the viscous device η_1 in Figure 8. The corresponding constitutive equation from Eq.2 above is obtained as follows:

$$\frac{\eta}{E_1} \overset{\bullet}{\sigma} + \left(1 + \frac{E_0 e^{-\beta t}}{E_1}\right) \sigma = \eta \overset{\bullet}{\varepsilon} + E_0 e^{-\beta t} \varepsilon$$
(3)

Assuming that, at t=0, an instantaneous stress $\sigma = \sigma_0$ is applied, and, with substituting the boundary condition $\varepsilon_{(0)} = 0$ and solving the above differential equation, we obtain the creep equation of the revised generalized Kelvin model as follows:

$$\varepsilon_{(t)} = \frac{\sigma_0}{E_1} + \frac{\sigma_0}{E_0 e^{-\beta t}} \left[1 - \exp\left(-\frac{E_0 e^{-\beta t}}{\eta} t\right) \right]$$
(4)

3.2 Nonlinear tensile viscoelastic-plastic creep model identification

3.2.1 Nonlinear viscoplastic body

From Figure 3, we can observe that, under the highest stress level, the strain of the rock is not limited to a particular value, but appears to be an accelerated flowing deformation; thus, when the tensile stress level exceeds the yield stress, the rock appears to be a viscoplastic material. Because the properties of a viscoplastic material are nonlinear, this material cannot be described by a traditional model like that given in Figure 6. Based on this, a nonlinear viscous model could be achieved through fitting and analog method, then combine it with a plastic device in parallel and to form a new nonlinear viscoplastic model by adopting the stress trigger method.

The accelerated creep portions of the curves in Figure 3 are fit by adopting the mixed power-exponent function given as follows:

$$\varepsilon = \varepsilon_0 + Ate^{nt} \tag{5}$$

Here, n is the coefficient of the accelerated creep exponential term that reflects the accelerated creep rate of the rock, where the rate of accelerated creep increases with increasing n.

Table 4. Fitting results of the accelerated creep portion for specimens I and II based on Eq.5.

	-1			
Specimen No.	${oldsymbol{\mathcal{E}}}_0$	A	п	R^2
Ι	6.6507E-4	9.2384E-30	0.2366	0.977
II	8.6849E-4	3.3742E-18	0.0697	0.985

From the values of \mathbb{R}^2 listed in Table 4, we can observe that the fitting results to the direct tensile accelerated creep portion of the curves in Figure 3 using Eq.5 are relatively good. For the two-component viscoplastic model, when $\sigma > \sigma_s$, where σ_s is the yield strength which is the stress required to cause the local macroscopic plastic deformation, the creep equation is as follows.

$$\varepsilon = \frac{\sigma - \sigma_s}{\eta} t + \varepsilon_0 \tag{6}$$

Setting Eq.5 and 6 equal to one another and taking $\frac{\sigma - \sigma_s}{\eta_0} = A$ the nonlinear equation of η for the viscous device can be obtained as follows.

$$\eta(n,t) = \frac{\eta_0}{e^{nt/t_0}} = \frac{\eta_0}{e^{nt}}$$
(7)

Here, t_0 is the reference time to be set as 1 and η_0 is the initial viscosity coefficient of accelerated creep.



Fig. 9. Nonlinear viscoplastic model

The obtained nonlinear viscoplastic model is as shown in Figure 9, the creep equation of which is given as follows.

$$\varepsilon_{(t)} = \frac{H(\sigma_0 - \sigma_s)}{\eta(n, t)} = \frac{H(\sigma_0 - \sigma_s)}{\eta_0} t \exp(nt)$$
(8)

Here, H is given by the following.

$$H(\boldsymbol{\sigma}_{0}-\boldsymbol{\sigma}_{S}) = \begin{cases} 0 & (\boldsymbol{\sigma}_{0} \leq \boldsymbol{\sigma}_{S}) \\ \boldsymbol{\sigma}_{0}-\boldsymbol{\sigma}_{S} & (\boldsymbol{\sigma}_{0} \leq \boldsymbol{\sigma}_{S}) \end{cases}$$

3.2.2 Nonlinear viscoelastic-plastic creep model

A series combination of the revised generalized Kelvin model and the newly proposed nonlinear viscoplastic model is presented to obtain a new nonlinear viscoelastic-plastic creep model appropriate for rock, as shown in Figure 10.



Fig. 10. Nonlinear viscoelastic-plastic creep model.

When $\sigma_0 \le \sigma_s$, the model is transformed into the revised generalized Kelvin model. When $\sigma_0 > \sigma_s$, the model is a five-component nonlinear viscoelastic-plastic model. The corresponding state equation is as follows.

$$\sigma = \sigma_1 = \sigma_2 = \sigma_3, \quad \sigma_1 = E_1 \varepsilon_1, \quad \sigma_2 = E_2 \varepsilon_2 + \eta_2 \varepsilon_2$$

$$\sigma_3 = \sigma_s + \eta_0 \varepsilon_3 / t e^{nt}, \quad \varepsilon = \varepsilon_1 = \varepsilon_2 = \varepsilon_3, \quad E_2 = E_0 e^{-\beta t}$$
(9)

The above formula is arranged to obtain the corresponding constitutive equation as follows.

$$E_{1}E_{0}e^{-\beta t}\varepsilon + E_{1}\eta_{2}\varepsilon = (E_{1} + E_{0}e^{-\beta t})\sigma + \frac{(\eta_{0} + E_{1}te^{nt})\eta_{2}\sigma}{\eta_{0}} + \frac{E_{1}(E_{0}e^{-\beta t} + \eta_{2}n)(\sigma - \sigma_{s})te^{nt}}{\eta_{0}}$$
(10)

In the above equation, σ and ε are the total stress and total strain of the model, respectively; σ_1 , σ_2 and σ_3 are respectively the stresses corresponding to sections 1, 2, and 3 in Figure 10; ε_1 , ε_2 , and ε_3 are respectively the strains corresponding to sections 1, 2, and 3 in Figure 10; E_0 , E_1 and E_2 are the elasticity parameters of the rock materials; and η_2 and η_0 are the viscosity parameters of the rock, wherein η_0 is the initial value of $\eta_{(n,t)}$ shown in Figure 10.

Under the effect of constant stress, we assume $\sigma = \sigma_0$, which is substituted into Eq.10, and the differential equation is solved to obtain the following creep equation of the five-

device nonlinear viscoelastic-plastic model given in Figure 10.

$$\varepsilon_{(t)} = \frac{\sigma_0}{E_1} + \frac{\sigma_0}{E_0 e^{-\beta t}} \left[1 - \exp\left(-\frac{E_0 e^{-\beta t}}{\eta_1} t\right) \right] + \frac{\sigma_0 - \sigma_s}{\eta_0} t \exp(nt) \quad (11)$$

3.2.3 Determination of the model parameters

Before determining the model parameters, the Boltzmann superposition principle is first utilized to transform the creep curves under graduated loading conditions given in Figure3 into those under the respective loading conditions shown in Figure 11.

Based on the test results, the present study utilizes the 1stOpt optimization software to identify the model parameters. The core of the software is the Universal Global Optimization, the greatest feature of which that it alleviates the problems associated with establishing suitable initial values when using an iteration method in the optimization. As such, the user need not determine initial parameter values, which are determined randomly by the software, and the optimal solution is obtained through its unique global optimizing calculation.

When $\sigma_0 \leq \sigma_s$, Eq.4 of the revised generalized Kelvin model is adopted to conduct fitting to the test curve, whereas, for $\sigma_0 > \sigma_s$, Eq.10 of the nonlinear viscoelastic-plastic model is adopted to conduct fitting to the test curve. The related parameters obtained are listed in Table 5. From the R² values of the fitting results, which are all greater than 0.98, we observe that the model fits the data quite well.

Specimen No.	Stress (MPa)/ stress-strength ratios	E ₀ (GPa)	E ₁ (GPa)	β (1)	η_{1} (GPa•h)	η (GPa•h)	<i>n</i> (1)	R^2
	0.44 /0.25	31.9525	3.7221	0.0102	155.9995	-	-	0.9986
Ι	0.74 /0.40	24.7676	3.4996	0.0157	121.1304	-	-	0.9961
	1.04 /0.58	19.5424	2.9551	0.0147	201.4115	-	-	0.9971
	1.34 /0.75	11.5821	2.6397	0.0126	119.5005	9.4728E11	0.5838	0.9989
	0.53 /0.30	22.7819	3.1443	0.0054	147.3809	-	-	0.9870
II	0.84 /0.47	25.7224	2.3914	0.0062	336.9538	-	-	0.9979
	1.14 /0.64	20.7722	2.1809	0.0044	450.3242	-	-	0.9988
	1.45 /0.81	7.5487	1.2150	0.0996	52.4185	3.8838 E10	0.0283	0.9999

Table 5. Fitted parameters of the nonlinear viscoelastic-plastic creep model.

Figure 11 presents comparisons between the creep test curves and the curves derived from the theoretical model for fine sandstone under conditions of tensile stress. From the results, we observe that the two curves are rather identical, which not only demonstrates the optimization capability of the 1stOpt software, but also shows that the revised generalized Kelvin model and the nonlinear viscoelasticplastic model proposed in the present study can accurately reflect the properties of attenuation creep, steady creep, and accelerated creep for fine sandstone under the effect of tensile stress, verifying the correctness of the proposed model.





Fig. 11. Comparison between the theoretical curves and test curves. (a) specimen I. (b) specimen II

4. Conclusions

The tensile creep and compression creep for fine sandstone exhibit similar curve shapes; however, the ratios of the creep strain to the total strain in the uniaxial tensile test are obviously higher than those in the compression test, indicating that the creep ability of rock in the uniaxial tensile process is higher than that in the uniaxial compression process.

A new revised generalized Kelvin model was proposed in the present study based on the elasticity modulus in the approximately linear portions of the isochronous stressstrain curves of the tensile creep, and the corresponding creep equation for this model was derived. By adjusting the value of the nonlinear constant β in the model, it is possible to describe the properties of attenuation creep and steady creep for different rock materials under stress conditions lower than the yield stress.

The revised generalized Kelvin model was combined in series with the proposed nonlinear viscoplastic model to form a new nonlinear viscoelastic-plastic creep model to describe the attenuation creep, steady creep, and the accelerated creep properties of rock materials.

Based on the creep test results, the 1stOpt optimization software was used to identify the model parameters, and the squared correlation coefficients obtained in the fitting process were all greater than 0.98. Comparison of the test curves with the model curves indicated that they were nearly identical, verifying that the proposed creep model can accurately describe the creep properties of fine sandstone under tensile stress.

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